

Solutions and Grading Key, Creative Response

Problem 1

(a, 8)

The force F_o must accelerate N blocks each of mass m :

$$F_o = Nma, \text{ so that } a = F_o / Nm. \quad (4)$$

Let T be the tension dragging $N-n$ blocks to the right. $F = ma$ gives

$$T = (N-n)ma \quad (2)$$

Using a ,

$$T = (N-n) F_o / N = F_o (1 - n/N). \quad (2)$$

(b, 12)

Each block has exerted on it the friction force μmg . (2)

$F = ma$ applied to the entire N -block system gives

$$F_o - N\mu mg = Nma \quad (2)$$

Solving for acceleration,

$$a = (F_o / Nm) - \mu g \quad (2)$$

Let T be the tension dragging $N-n$ blocks to the right. $F = ma$ gives

$$T - (N-n)\mu mg = (N-n)ma \quad (2)$$

so that upon substituting for a ,

$$T = (N-n)m [\mu g + (F_o / Nm) - \mu g] \quad (2)$$

which gives the final answer,

$$T = F_o (1 - n/N), \text{ the same as without friction.} \quad (2)$$

Problem 2

(a, 10) Work energy theorem: $W = \Delta K$ (2)

Setting up the work energy theorem,

$$[mg \sin \theta - \mu N]L = \frac{1}{2}mv^2 - 0, \quad (4)$$

where N is the normal force. But from $F = ma$ normal to the roof, $N = \mu mg \cos \theta$

(2)

thus, $L = \frac{1}{2}v^2 / [g(\sin \theta - \mu \cos \theta)]$

and numerically,

$$L = \frac{1}{2}(3.5 \text{ m/s})^2 / [(9.8 \text{ m/s}^2)(\sin 30^\circ - \frac{1}{4} \cos 30^\circ)] = 2.2 \text{ m}. \quad (2)$$

(-1 for incorrect units; -1 for gross neglect of significant figures)

Alternative method:

$$F = ma \text{ tangent to the roof gives } a = g \sin \theta - \mu N \quad [4]$$

$$= g \sin \theta - \mu g \cos \theta \quad [2]$$

$$= 2.8 \text{ m/s}^2. \quad [1]$$

$$\text{Using } v^2 - v_0^2 = 2a(x - x_0) \quad [2]$$

$$\text{leads to } L = v^2 / 2a = 2.2 \text{ m}. \quad [1]$$

(b, 10) Projectile problem:

$$v_{0y} = (3.5 \text{ m/s}) \sin 30^\circ = 1.75 \text{ m/s}, \quad (1)$$

$$v_{0x} = (3.5 \text{ m/s}) \cos 30^\circ = 3.03 \text{ m/s} = 3.0 \text{ m/s}. \quad (1)$$

For the x -direction,

$$x = v_{0x} t, \quad (2)$$

For the y -direction (taking down positive)

$$y = v_{0y} t + \frac{1}{2}gt^2. \quad (2)$$

Solve the equations simultaneously by some appropriate strategy. For instance, eliminating t in the y -equation by using

the x -equation: $y = v_{0y} (x/v_{0x}) + \frac{1}{2}g(x/v_{0x})^2$ (2)

Putting in numerical values, solving the quadratic equation correctly, choosing the + root in the quadratic formula (with the - root one gets an $x < 0$, which is *inside* the house), one finds $x = 1.9 \text{ m}$.

(2)

Yes, the box lands in the flowers.

Problem 3

(a, 3)

By conservation of energy,

$$\frac{1}{2}I\omega_0^2 + mgh = K_f + mgR$$

so that

$$K_f = \frac{1}{2}I\omega_0^2 + mg(h - R) \quad (3)$$

(b, 17)

Let (v_x, v_y) = velocity of the ball immediately *after* the collision is over.

Given:

$$\beta K_f = \frac{1}{2}m(v_x^2 + v_y^2) \quad (2)$$

Since β and K_f are known, if either of v_x or v_y is known then so is the other. Let us find v_x .

(Note that linear momentum is *not* conserved.)

In the collision, the angular frequency goes from ω_0 to θ .

Let f be the average force of friction that acts horizontally on the ball during the collision. Then from torque = $\Delta L/\Delta t$,

$$fR = I\omega_0/\Delta t$$

where Δt is the collision time. (5)

Now apply $F = m \Delta v/\Delta t$ in the x -direction.

The x -component of velocity changes

from θ to v_x during time Δt , and the horizontal force is f .

$$f = m v_x/\Delta t. \quad (5)$$

From the torque equation we have $f\Delta t = I\omega_0/R$, so combining the torque and force equations,

$$I\omega_0/R = m v_x \quad (2)$$

and solving for v_x ,

$$v_x = I\omega_0/mR \quad (2)$$

Now v_x is known. Thus, so is v_y ,

$$v_y = [2\beta K_f/m - v_x^2]^{1/2} \quad (1)$$

Problem 4

(a, 7) From $F = ma$, since the orbit is circular so that the acceleration is v^2/R ,

$$(1 + \Gamma)(GMm/R^2) = mv^2/R \quad (2)$$

Using $v = 2\pi R/T$ where T is the period, (2)

$F = ma$ becomes (after some re-arrangement)

$$T^2 = (4\pi^2/GM)R^3(1 + \Gamma)^{-1}$$

Notice that $T_0^2 = (4\pi^2/GM)R^3$, where T_0 is the period of the planet's orbit with purely the Newtonian force. Hence, $T = T_0(1 + \Gamma)^{-1/2}$ (2)

Since $\Gamma \ll 1$, by the binomial expansion

$$T \approx T_0(1 - \frac{1}{2}\Gamma) \quad (1)$$

(b, 7)

In elapsed time

$$\Delta t = T_0, \text{ the planet turns through the angle } 2\pi + \delta \quad (2)$$

This angle can also be equated to ωT_0 where ω is the planet's orbital angular velocity, viz.,

$$\begin{aligned} \omega &= 2\pi / T \\ &= (2\pi / T_0)(1 - \frac{1}{2}\Gamma)^{-1} \end{aligned} \quad (2)$$

by virtue of the binomial expansion again,

$$\omega \approx (2\pi / T_0)(1 + \frac{1}{2}\Gamma) \quad (1)$$

Therefore, we complete the calculation:

$$\begin{aligned} 2\pi + \delta &= \omega T_0 \\ &= (2\pi)(1 + \frac{1}{2}\Gamma) \end{aligned}$$

cancelling the 2π ,

$$\delta = \pi\Gamma. \quad (2)$$

Problem 4, continued

(c, 2)

$$\begin{aligned}\delta &= \pi \Gamma \\ &= 6\pi v^2/c^2\end{aligned}$$

From $F = ma$ above, we may approximate v^2 with GM/R , (keeping the Γ from $F = ma$ would give δ here to second order in Γ), so that

$$\delta = 6\pi GM/Rc^2 \quad (2)$$

(d, 4)

For one lap around the Sun, from the given numerical data,

$$\begin{aligned}\delta_l &= 6\pi GM/Rc^2 \\ &= 6\pi (6.7 \times 10^{-11} \text{ Nm}^2/\text{kg}^2)(2.0 \times 10^{30} \text{ kg})(3.0 \times 10^8 \text{ m/s})^{-2}(5.8 \times 10^{10} \text{ m})^{-1} \\ &= 4.8 \times 10^{-7} \text{ radians}\end{aligned} \quad (2)$$

In 100 years, Mercury makes

$$(365 / 88) \times 100 = 415 \text{ laps around the Sun.} \quad (1)$$

For 415 laps,

$$\begin{aligned}\delta_{415} &= 415 (4.8 \times 10^{-7} \text{ radians}) \\ &= 2 \times 10^{-4} \text{ radians} \\ &= (2 \times 10^{-4})(180/\pi) \text{ degrees} \\ &= (2 \times 10^{-4})(180/\pi)(3600'') \\ &= 41''.\end{aligned} \quad (1)$$